

Home Search Collections Journals About Contact us My IOPscience

Shock formation and traffic jam induced by a crossing in the 1D asymmetric exclusion model

This article has been downloaded from IOPscience. Please scroll down to see the full text article. 1993 J. Phys. A: Math. Gen. 26 6625 (http://iopscience.iop.org/0305-4470/26/23/013)

View the table of contents for this issue, or go to the journal homepage for more

Download details: IP Address: 171.66.16.68 The article was downloaded on 01/06/2010 at 20:10

Please note that terms and conditions apply.

# Shock formation and traffic jam induced by a crossing in the 1D asymmetric exclusion model

Takashi Nagatani

College of Engineering, Shizuoka University, Hamamatsu 432, Japan

Received 26 July 1993

Abstract. The asymmetric simple-exclusion model on two-crossing one-dimensional lattices is considered for studying shock formation and traffic jam induced by a crossing. The condition for shock formation is derived for car densities  $p_1$  and  $p_2$  of the first and second street. The phase diagram and the dependence of the traffic current on the car densities are shown. We compare the present result with that of a car accident. Also, we investigate the shock formation induced by both a crossing and a car accident. It is found that the dynamical transitions occur successively with increasing car density. The phase diagram is shown.

### 1. Introduction

The one-dimensional (1D) exclusion model is one of the simplest examples of a driven system [1, 2]. The model has been extensively studied for understanding systems of interacting particles [3, 4]. The 1D exclusion models are used to study the microscopic structure of shocks [5, 6] and are closely linked to growth processes [7-9]. The 1D asymmetric simple-exclusion model can be formulated as traffic jam problems.

Recently, traffic problem have attracted considerable attention. Traffic simulations based on various hydrodynamic models have provided much insight [10, 11]. However, the simulation of traffic flow is a formidable task since it involves many degrees of freedom. Cellular automaton (CA) models are being applied successfully to simulations of complex physical systems [12, 13]. Very recently, Biham *et al* [14] applied the two-dimensional asymmetric exclusion model to the traffic-jam problem. They found that a dynamical jamming transition between the low-density moving phase and the high-density jamming phase occurs at a critical density of cars. Also, in order to simulate freeway traffic, Nagel and Schreckenberg [15] extended the 1D asymmetric exclusion model to take into account car velocity. They showed that a transition from laminar traffic flow to start-stop waves occurs with increasing car density as is observed in real freeway traffic.

In real traffic-flow systems, the traffic jam is frequently induced by a single crossing when a street (first street) crosses with another street (second street). The congested street prevents cars from crossing its street. As soon as a single street begins to be congested, a traffic jam spreads from the crossing throughout another street. The occurence of a traffic jam on the first street strongly depends on whether or not the second street (crossing with the first street) becomes congested. The traffic jam appears as a shock (a discontinuity of densities) which separates between the low-density and the high-density traffic flow. The mechanism of shock formation in the 1D asymmetric simple-exclusion model was examined by Janowsky and Lebowitz [5]. They studied a shock structure when the translation invariance is broken by the insertion of a blockage. The model corresponds to the traffic-jam problem induced by a car accident. The

mechanism of shock formation induced by the car accident will be different from that induced by the crossing. However, the shock formation and the traffic jam induced by a crossing have not been studied until now.

In this paper, we present the extended verision of the 1D asymmetric simple-exclusion model in which two one-dimensional lattices cross one another. We investigate the shock formation and traffic jam induced by a crossing. We find the phase diagram for shock formation against the car densities. We show the dependence of traffic current on the car densities. We compare our result with that of Janowsky and Lebowitz [5]. We find that the shock formation induced by a crossing is different from that of a car accident. Furthermore, we study the combined effect of both crossing and car accident on shock formation and traffic jam.

The organization of the paper is as follows. In section 2, we present our model. The traffic current and the phase diagram for shock formation are derived. The result is compared with that of Janowsky and Lebowitz. In section 3, we study the shock formation process induced by both crossing and car accident. We show the phase diagram for shock formation. Section 4 contains a brief summary.

# 2. Model and simulation result

Our model is defined on two one-dimensional lattices in which each lattice consists of L sites with periodic boundary conditions and the first lattice crosses with the second. Each site is either occupied by one car or is empty. Figure 1 shows the schematic representation of our model. The arrow pointing up (right) represents the car moving up (right). For an arbitrary configuration, one update of the system consists of the two steps. The first step is performed in parallel for all the cars on the first street (lattice) and the second step is performed in parallel for all the cars on the second street (lattice). The move or stop of cars in each step is the same as the 1D asymmetric simple-exclusion model. In each lattice, each arrow moves forward one step unless the forward nearest-neighbour site is occupied by another arrow. If an arrow is blocked ahead by another arrow, it does not move even if the blocking arrow moves out of the site during the same time step. In this model, the total number of cars on each street is conserved. The traffic problem on two streets connected by a crossing is reduced to its simplest form. The essential features are maintained. These features include the simultaneous flow of cars on two streets which cannot overlap and a car on the first (second) street at the crossing prevents cars on the second (first) street from going ahead. In the limit of no cars on the first or second street, our model reproduces the 1D asymmetric simple-exclusion model.

We have performed simulations of the CA model starting with an ensemble of random initial conditions where the system size L = 2000, and the initial densities of cars on the first and second streets are  $p_1 = 0.0-1.0$  and  $p_2 = 0.0-1.0$ . Each run is calculated up to 10000 time steps. The data are averaged over 50 runs. We present the simulation result obtained by the procedure explained above. We consider the traffic flow under the condition of a constant car density  $p_2$  on the second street. Figure 2 shows the plot of the mean traffic-current  $J_1$  on the first street against the car density  $p_1$  on the first street for  $p_2 \leq 0.5$  (circles),  $p_2 = 0.65$  (triangles) and  $p_2 = 0.8$  (squares). The mean current J is obtained by averaging over 5000 time steps except for the initial stage. For density  $p_2$  lower than 0.5, the traffic flow of the second street does not affect the traffic current  $J_1$  on the first street is independent of that on the second street. For  $p_2 = 0.65$ , the traffic flow on the first street is independent of that on the second street. For  $p_2 = 0.65$ , the traffic current  $J_1$  becomes saturated at the car density  $p_1 = 0.35 \pm 0.01$ , remains at the constant value  $J_1 = 0.35 \pm 0.01$  until  $p_1 = 0.66 \pm 0.01$  and then decreases linearly with  $p_1$ . Similarly, for  $p_2 = 0.8$ , the traffic current  $J_1$  becomes





Figure 2. The plot of the mean traffic-current  $J_1$  on the first street against the car density  $p_1$  on the first street for  $p_2 \leq 0.5$  (circles),  $p_2 = 0.65$  (triangles) and  $p_2 = 0.8$  (squares).

saturated at the density  $p_1 = 0.2 \pm 0.01$ , remains at the constant value  $J_1 = 0.2 \pm 0.01$  until  $p_1 = 0.81 \pm 0.01$  and then decreases linearly with  $p_1$ . The saturation of traffic current  $J_1$  is due to the shock formation on the first street. The shock is a discontinuity of density which separates between the low density  $p_{1,low}$  and the high density  $p_{1,lowh}$ . For illustration, figure 3 shows the typical configuration of cars up to 500 time steps for  $p_1 = 0.5$  and  $p_2 = 0.8$  where the system size L = 190. The upper and lower configurations represent respectively those on the second and first streets. The vertical direction indicates that of moving cars. The horizontal direction indicates that of time. A car is indicated by a dot. The trajectory of a car is represented by a line. The crossing is on the tops of the first and second streets. The shock (discontinuity of densities) is positioned near the centre of the first street. Figure 4 shows the density profiles  $p_1(x)$  and  $p_2(x)$  of the first and second streets against the distance x for the initial densities  $p_1 = 0.5$  and  $p_2 = 0.8$  where the system size L = 2000 and the density is averaged over 5000 time steps. The shock appears at x = 1050 on the first street. The shock separates between the low density  $p_{1,\text{low}} = 0.2 \pm 0.01$  and the high density  $p_{1,\text{bigh}} = 0.8 \pm 0.01$ . The high density  $p_{1,\text{high}}$  agrees with the density  $p_2$  of the second street. When the shock formation occurs, the high density  $p_{1,\text{high}}$  of the shock is determined by the density  $p_2$ . We find the following



Figure 3. The typical configuration of cars up to 500 time steps for  $p_1 = 0.5$  and  $p_2 = 0.8$  where the system size L = 190. The upper and lower configurations represent, respectively, those on the second and first streets. The crossing is on the tops of the first and second streets. The upper direction indicates that of moving cars. The right direction indicates that of time. A car is indicated by a dot. The trajectory of a car is represented by a line. A shock (discontinuity of densities) appears near the centre of the first street.



Figure 4. The density profiles  $p_1(x)$  and  $p_2(x)$  of the first and second streets against the distance x for the densities  $p_1 = 0.5$  and  $p_2 = 0.8$  where the system size L = 2000 and the density is averaged over 5000 time steps. The crossing is positioned at x = 2000. The shock appears at x = 1050 on the first street. The shock separates between the low density  $p_{1,\text{low}} = 0.2 \pm 0.01$  and the high density  $p_{1,\text{high}} = 0.8 \pm 0.01$ . The high density  $p_{1,\text{high}}$  agrees with the density  $p_2$  of the second street.

relationships between the low and high densities:

$$p_{1,\text{high}} = p_2 \quad \text{and} \quad p_{1,\text{low}} + p_{1,\text{high}} = 1.$$
 (1)

The second equality can be derived from the conservation law of traffic current after and before





the shock on the first street. The mean velocity  $\langle v \rangle$  in the 1D asymmetric exclusion model is given by  $\langle v \rangle = 1$  for  $p \leq 0.5$  and  $\langle v \rangle = (1 - p)/p$  for p > 0.5. The conservation law of current after and before the shock is given by  $p_{1,\text{low}}\langle v_{1,\text{low}}\rangle = p_{1,\text{bigh}}\langle v_{1,\text{high}}\rangle$  where  $\langle v_{1,\text{low}}\rangle$  and  $\langle v_{1,\text{high}}\rangle$  are the mean velocities before and after the shock. They are given by  $\langle v_{1,\text{low}}\rangle = 1$  and  $\langle v_{1,\text{high}}\rangle = (1 - p_{1,\text{high}})/p_{1,\text{high}}$ . Therefore, the relationship (1) is satisfied. The shock appears when the density  $p_2 \geq 0.5$ . The shock front shifts backward with the initial density  $p_1$ . The position  $x_s/L$  of the shock is determined by the initial density  $p_1$ . It is given by

$$x_{\rm s}/L = (p_{1,\rm high} - p_1)/(p_{1,\rm high} - p_{1,\rm low})$$
 for  $p_1 \ge p_{1,\rm low}$ . (2)

Figure 5 shows the phase diagram between  $p_1$  and  $p_2$ . The shock does not appear under the following condition:

$$p_1 < 1 - p_2$$
. (3)

A shock appears only on the first street under the condition

$$p_1 > 1 - p_2$$
 and  $p_1 < p_2$ . (4)

The densities  $p_{1,\text{high}}$  and  $p_{1,\text{low}}$  in the front and back of the shock are given by (1). The shock position is given by (2). A shock appears only on the second street under the condition

$$p_1 > 1 - p_2$$
 and  $p_1 > p_2$ . (5)

The densities  $p_{2,high}$  and  $p_{2,low}$  in the front and back of the shock are given by the same equation as (1) where the subscript 1 (2) is replaced by 2 (1). The shock position is also given by the same as (2) where the subscript 1 is replaced by 2. Shock appears on the street with lower car density. Shocks never appear simultaneously on both streets.

We compare our result with that of a car accident. Janowsky and Lebowitz investigated the shock formation induced by a car accident in the 1D asymmetric simple-exclusion model [5]. They showed the mechanism of shock formation when the translation invariance is broken by the insertion of a blockage: this reduces the rates at which particles jump across the bond by a factor r (0 < r < 1). Then, shock formation occurs when the following condition is satisfied:

$$r/(r+1) 
(6)$$

6630 T Nagatani



Figure 6. The plot of traffic currents  $J_1$  (squares) and  $J_2$  (circles) against the car density  $p_1$  on the first street for  $p_2 = 0.4$  and  $r_1 = 0.35$ .

where p is the car density and r is the rate representing the degree of blocking. The shock separates the low density  $p_{\text{low}}$  from the high density  $p_{\text{high}}$ . The densities  $p_{\text{low}}$  and  $p_{\text{high}}$  are given by

$$p_{\text{low}} = r/(r+1)$$
 and  $p_{\text{high}} = 1/(r+1)$ . (7)

The condition (6) of shock formation induced by a blockage is definitely different from equation (4) and (5) of shock formation induced by a crossing. The densities  $p_{\text{high}}$  and  $p_{\text{low}}$  of (7) in the front and back of the shock are also different from (1) by a crossing.

Our model is fully deterministic since updating is performed in parallel. Even if updating is performed on a randomly chosen site, the results are the same.

### 3. Combined effect of both crossing and car accident

We consider shock formation and traffic jam induced by both the crossing and the car accident. We derive the phase diagram for shock formation and traffic jam. We assume that a car accident occurs at a position  $x_a$  on the first street. The distance  $x_a$  is measured from the position at which cars pass just the crossing: x = 0 and x = L are positioned just after and just before the crossing. When a car flows through the car accident, it passes with probability  $r_1$  or it stops with probability  $1 - r_1$ . The blockage of the car accident reduces the rate at which cars move across the position of car accident by a factor  $r_1$  ( $0 < r_1 < 1$ ). We study the combined effect of both crossing and car accident on the shock formation and traffic jam. We put the blockage by car accident on the position  $x_a = L/2$  on the first street. We calculate the traffic currents  $J_1$  and  $J_2$  on the first and second streets. Figure 6 shows the plot of traffic currents  $J_1$  and  $J_2$ against the car density  $p_1$  on the first street for  $p_2 = 0.4$  and  $r_1 = 0.35$  where the system size L = 2000 and the currents are averaged over 5000 time steps. Cars on the first and second streets move at the maximum velocity  $v_{\text{max}} = 1.0$  until  $p_1 = 0.26 \pm 0.02$ . The traffic current  $J_1$  is proportional to the car density  $p_1$  on the first street. The traffic flow on the first street is not affected by that on the second street and the car accident on the first street. At density  $p_1$ higher than 0.26, a shock induced by the car accident appears on the first street. It propagates from the car accident throughout space with increasing  $p_1$ . The traffic current  $J_1$  on the first street becomes saturated and  $J_1 = 0.26 \pm 0.02$  when the shock on the first street is formed. When  $p_1$  becomes 0.5, the shock reaches the crossing. Then, the car density at the crossing becomes  $p_{high}$  (> 0.5). A new shock appears on the second street. Two shocks induced by the car accident and the crossing occur respectively on the first and second streets. The traffic current  $J_2$  on the second street decreases sharply at  $p_1 = 0.5$ . Both traffic currents  $J_1$  and  $J_2$ become the constant value  $J_1 = J_2 = 0.26 \pm 0.02$  until  $p_1 = 0.74 \pm 0.02$ . When the car density  $p_1$  becomes larger than 0.74, the shock on the first street disappears. Only the shock on the second street appears. We find the following:

 $J_2 = 0.4$ for  $0 < p_1 < 0.26$  $J_1 = p_1$ and (phase 1)  $J_2 = 0.4$  $0.26 < p_1 < 0.5$  $J_1 = 0.26$ for and (phase 2) (8)  $J_1 = J_2 = 0.26$ for  $0.5 \leq p_1 < 0.74$ (phase 3) for  $0.74 < p_1 < 1^{1}$  $J - 1 = J_2 = 1 - p$ (phase 5).

The dynamical phase transitions occur successively from the phase 1 (where the cars move with the maximum velocity), through phases 2 (the shock induced by the car accident appears on the first street) and 3 (two shocks induced by the car accident and the crossing appear respectively on the first and second streets), to phase 5 (a shock induced by the crossing appears only on the second street). Since the shock on the first street is induced by the car accident, the densities  $p_{1,\text{high}}$  and  $p_{1,\text{how}}$  in the front and back of shock on the first street are given by (7):

$$p_{1,\text{high}} = 1/(r_1 + 1) = 0.74$$
 and  $p_{1,\text{low}} = r_1/(r_1 + 1) = 0.26$  (9)

where  $r_1 = 0.35$ . Since the shock on the second street is induced by the crossing, the densities  $p_{2,\text{high}}$  and  $p_{2,\text{low}}$  in the front and back of shock on the second street are given by (1):

$$p_{2,\text{high}} = p_{1,\text{high}} = 0.74 \quad \text{and} \quad p_{2,\text{low}} = 1 - p_{2,\text{high}} = 0.26$$
for  $0.5 \leq p_1 < 0.74$ 

$$p_{2,\text{high}} = p_1 \quad \text{and} \quad p_{2,\text{low}} = 1 - p \quad \text{for} \quad 0.74 \leq -p_1 < 1.$$
(10)

For illustration, figure 7 shows the typical configurations of cars up to 500 time steps for  $p_1 = 0.7$ ,  $p_2 = 0.4$  and  $r_1 = 0.35$  where L = 190. The upper and lower configurations represent respectively those on the second and first streets. The upper direction indicates that of moving cars. The fight direction indicates that of time. A car is indicated by a dot. The trajectory of a car is represented by a line. The crossing is on the tops of the first and second streets. The car accident is positioned at the centre of the first street. The shocks (discontinuities of densities) appear on the first and second streets.

We calculate traffic currents in the case of higher car density  $p_2$ . Figure 8 shows the plot of traffic currents  $J_1$  and  $J_2$  against the car density  $p_1$  on the first street for  $p_2 = 0.85$  and  $r_1 = 0.35$  where the system size L = 2000 and the currents are averaged over 5000 time steps. Cars on the first and second streets move at the maximal velocity  $v_{\text{max}} = 1.0$  until  $p_1 = 0.15 \pm 0.02$ . The traffic current  $J_1$  is proportional to the car density  $p_1$  on the first street.

The traffic flow on the first street is not affected by that on the second street and the car accident on the first street. At density  $p_1$  higher than 0.15, a shock induced by the crossing appears on the first street. It propagates from the crossing throughout space with increasing  $p_1$ . The traffic current  $J_1$  on the first street becomes saturated and equals  $J_2$  ( $J_1 = J_2 = 0.15 \pm 0.02$ ) when the shock on the first street is formed. In this case of higher density  $p_2$ , the shock on the first street is induced by the crossing. In the previous case of lower density  $p_2$ , the shock on



Figure 7. The typical configurations of cars up to 500 time steps for  $p_1 = 0.7$ ,  $p_2 = 0.4$  and  $r_1 = 0.35$  where L = 190. The upper and lower configurations represent respectively those on the second and first streets. The upper direction indicates that of moving cars. The right direction indicates that of time. The crossing is on the tops of the first and second streets. The car accident is positioned at the centre of the first street. The shocks (discontinuities of densities) appear on the first and second streets.



6632

Figure 8. The plot of traffic currents  $J_1$  (squares) and  $J_2$  (circles) against the car density  $p_1$  on the first street for  $p_2 = 0.85$  and  $r_1 = 0.35$ .

the first street is induced by the car accident. Whether the shock on the first street is induced by the crossing or by the car accident depends on the car density  $p_2$  on the second street. When the car density  $p_1$  on the first street becomes larger than  $p_2$ , the shock on the first street disappears and a new shock on the second street is induced by the crossing. The current  $J_1$ becomes equal to  $J_2$  and decreases linearly with  $p_1$ . We find the following:

$J_1 = p_1$	and $J_2$	= 0.15	for	$0$	(phase 1)	
$J_1 = J_2 = 0.1$	5 for	0.15 ≼	$\leq p_1 < 0.85$	(phase 4)		(11)
$J_1 = J_2 = 1 -$	- p <sub>1</sub> fo	r 0.85	$5 \leq p_1 < 1$	(phase 5).		





The dynamical phase transitions occur successively from phase 1 (where the cars move with the maximum velocity), through phase 4 (the shock induced by the crossing appears on the first street), to phase 5 ( the shock induced by the crossing appears only on the second street). In this case of higher density  $p_2$ , the shock formation is not induced by the car accident but by the crossing. The dynamical phase transitions donot depend on the car accident. The transitions agree with those by the crossing explained in section 2. We calculate traffic flow for various values of  $p_2$ . We find the phase diagram for the shock formation process. Figure 9 shows the phase diagram between  $p_1$  and  $p_2$ . Six dynamical phases 1–6 appear. Each phase is indicated by a number. The shock does not appear in phases 1 and 6. In phase 2, a shock induced by the car accident appears on the first street. In phase 3, a shock induced by the car accident appears on the first street and a shock induced by the crossing appears on the first street. In phase 4, a shock induced by the crossing appears on the second street. The regions of each phase are divided by the following:

phase 1 for	$0 < p_1 \leqslant r_1/(1+r_1) \qquad \text{and} \qquad$	$p_1 < 1$	$- p_2$
phase 2 for	$r_1/(1+r_1) < p_1 < 1 - x_a/L$	and	$0 < p_2 \leqslant 1 - r_1/(1+r_1)$
phase 3 for	$1 - x_{\rm a}, /L < p_1 < 1/(1 + r_1)$	and	$0 < p_2 \leqslant 1 - r_1 / (1 + r_1) $
phase 4 for	$1-p_2 < p_1 \leqslant p_2$ and	$1 - r_1/(1$	$(12)$ $+r_1) < p_2 < 1$
phase 5 for	$p_1 > 1 - p_2, p_1 > p_2$ and	$p_1 > 1$	$1/(1+r_1)$
phase 6 for	$1/(1+r_1) < p_1 \leq 1-p_2$ .		

The densities before and after the shock induced by the crossing are obtained by (1). The densities before and after the shock induced by the car accident are given by (7).

#### 4. Summary

We presented the extended version of the ID asymmetric simple-exclusion model to investigate shock formation and traffic jam induced by a crossing. The extended model consists of twocrossing one-dimensional lattices. We calculated the traffic current on the system. We derived the condition of shock formation and showed the phase diagram. Also, we investigated the combined effect of crossing and car accident on shock formation. We showed that the dynamical phase transitions occur successively with increasing car density. We found the phase diagram.

## References

- [1] Derrida B, Evans M R, Hakim V and Pasquier V 1993 J. Phys. A: Math. Gen. 26 1493
- [2] Gwa L H and Spohn H 1992 Phys. Rev. Lett. 68 725
- [3] Liggett T M 1985 Interacting Particle Systems (New York: Springer)
- [4] Spohn H 1991 Large Scale Dynamics of Interacting Particles (Berlin: Springer)
- [5] Janowsky S A and Lebowitz J L 1992 Phys. Rev. A 45 618
- [6] Brainson M 1988 J. Stat. Phys. 51 863
- [7] Meakin P, Ramanlal P, Sander L M and Ball R C 1986 Phys. Rev. A 34 5091
- [8] Krug J and Spohn H 1988 Phys. Rev. A 38 4271
- [9] Family F and Vicsek T (eds) 1991 Dynamics of Fractal Surfaces (Singapore: World Scientific)
- [10] Gatnet N H and Wilson N H M (eds) 1987 Transportation and Traffic Theory (New York: Elsevier)
- [11] Leutzbach W 1988 Introduction to the Theory of Traffic Flow (Berlin: Springer)
- [12] Wolfram S 1986 Theory and Applications of Cellular Automata (Singapore: World Scientific)
- [13] Kaneko K 1990 Formation, Dynamics and Statistics of Patterns vol 1, ed K Kawasaki, M Suzuki and A Onuki (Singapore: World Scientific) p 1
- [14] Biham O, Middleton A A and Levine D 1992 Phys. Rev. A 46 R6124
- [15] Nagel K and Schreckenberg M 1992 J. Physique I 2 2221